

Lecture 6 summary

2. Classical statistics in non-equilibrium

2.1 Master equation

- for a discrete state  $n$  (particle number)

$$\frac{d}{dt} p_n(t) = \sum_{\substack{m \\ m \neq n}} [W_{nm} p_m(t) - W_{mn} p_n(t)]$$

$n < m$                        $m < n$   
 gain                                  loss

- for one-step processes (birth-death):  $W_{nn'} = r_{n'} \delta_{n, n'-1} + g_{n'} \delta_{n, n'+1}$   
recomb.

$$\dot{p}_n = r_{n+1} p_{n+1} + g_{n-1} p_{n-1} - (r_n + g_n) p_n$$

$p_n$  - probability of finding the system in state  $n$

$g_n$  - generative rate (births)

$r_n$  - recombination rate (deaths)

Classification

- (i)  $r_n, g_n = \text{const}$  - random walk
- (ii)  $r_n, g_n$  linear in  $n$  - random process
- (iii)  $r_n, g_n$  nonlinear in  $n$  - bimolecular recomb.

Special case of (i):  $g_n = g, r_n = 0$  - Poisson process

Stationary solution of master equation  $p_n^*$

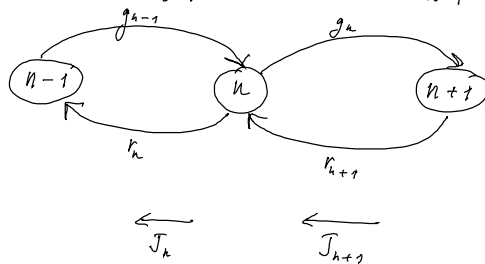
We can write the equation for the stationary solution  $p_n^*$  as

(1)  $0 = J_{n+1} - J_n$  with

$$J_n = r_n p_n^* - g_{n-1} p_{n-1}^*$$

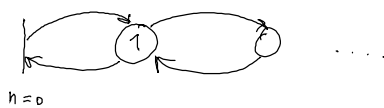
$J_{n+1}$  - incoming probability flow in state  $n$ ;  $n+1 \rightarrow n$

$J_n$  - outgoing probability flow in state  $n$ ;  $n \rightarrow n-1$



Boundary conditions  $n=0$ :

$$r_n = 0, p_{n=0} = 0$$



(no probability of an individual dying if there are none present;  
 $n$  is non-negative integer since we cannot have a negative number of individuals)

$$J_0 = 0 \quad (J_0 = r_0 p_0^* - g_{0-1} p_{0-1}^* = 0)$$

We now sum (1) so

$$0 = \sum_{n'=0}^{n-1} (J_{n'+1} - J_{n'}) = J_n - \underbrace{J_0}_{=0}$$

$$\Rightarrow J_n = 0 \Rightarrow p_n^* = \frac{g_{n-1}}{r_n} p_{n-1}^* \quad (*)$$

$$p_n^* = p_0^* \prod_{n'=1}^n \frac{g_{n'-1}}{r_{n'}} \quad \text{stationary solution}$$

### Detailed balance interpretation

The condition  $J_n = 0$  can be seen as a detailed balance requirement.

### Rate equations

We notice that the mean of  $N$  satisfies

$$\begin{aligned} \frac{d}{dt} \langle N \rangle &= \sum_{n=0}^{\infty} n \dot{p}_n = \sum_{n=0}^{\infty} n (r_{n+1} \underbrace{p_{n+1}}_{\tilde{n}} - r_n p_n) + \sum_{n=0}^{\infty} n (g_{n-1} \underbrace{p_{n-1}}_{\tilde{n}} - g_n p_n) = \\ &= \sum_{\tilde{n}=n}^{\infty} (\tilde{n} - 1) r_{\tilde{n}} p_{\tilde{n}} - \sum_{n=0}^{\infty} n r_n p_n + \sum_{\tilde{n}=n}^{\infty} (\tilde{n} + 1) g_{\tilde{n}} p_{\tilde{n}} - \sum_{n=0}^{\infty} n g_n p_n = \\ &\quad \text{0 since } r_0 = 0 \quad \text{0 since } \tilde{n} + 1 = 0 \text{ for } \tilde{n} = -1 \\ &= \sum_{n=0}^{\infty} g_n p_n - \sum_{n=0}^{\infty} r_n p_n \end{aligned}$$

$$\langle \dot{N} \rangle = \langle g_n \rangle - \langle r_n \rangle$$

NB: for nonlinear processes this does not give a closed equation for mean values since

$$\langle N^2 \rangle \neq \langle N \rangle^2, \text{ but a hierarchy for moments } \frac{d}{dt} \langle N^k \rangle.$$

The corresponding deterministic equation is that

which would be obtained by neglecting fluctuations:

$$\frac{dN}{dt} = g_n - r_n \quad \text{Notice that a stationary state occurs deterministically when (no evolution w.r. time):}$$

$$g_n = r_n \quad \text{Corresponding to this, notice that the maximum value of } p_n^* \text{ occurs when:}$$

$$\frac{p_n^*}{p_{n-1}^*} \simeq 1, \text{ which from } (*) \text{ corresponds to}$$

$$g_{n-1} = r_n$$

For sufficiently large  $N$ ,  $g_n = r_n$  and  $g_{n-1} = r_n$  are essentially the same. Thus, the modal value of  $N$ , which corresponds to  $g_{n-1} = r_n$ , is the stationary stochastic analogue of the deterministic steady state that corresponds to  $g_n = r_n$ .

Example: chemical reaction  $X \rightleftharpoons A$

We treat the case of a reaction  $X \xrightleftharpoons[k_2]{k_1} A$  in which it is assumed that  $A$  is a fixed concentration  $a$  ( $X$  plays a role of random variable  $N(t)$ ). Therefore:

$$g_n = k_2 a$$

$$r_n = k_1 n$$

The Master equation takes form:

$$\dot{p}_n = k_2 a p_{n-1} + k_1 (n+1) p_{n+1} - (k_1 n + k_2 a) p_n$$

generating function

To solve the equation, we introduce the generating function

$$G(s, t) = \sum_{n=0}^{\infty} s^n p_n(t) \quad \text{so that}$$

$$\frac{\partial}{\partial t} G(s, t) = k_2 a (s-1) G(s, t) - k_1 (s-1) \frac{\partial}{\partial s} G(s, t) \quad (* *)$$

If we substitute  $q_0(s, t) = G(s, t) e^{\frac{-k_2 a s}{k_1}}$  the equation  $(* *)$  becomes

$$\frac{\partial}{\partial t} q_0(s, t) = k_1 (s-1) \frac{\partial}{\partial s} q_0(s, t) \quad \left| \begin{array}{l} \frac{\partial}{\partial t} = \frac{\partial}{\partial t} \\ \exp = e \end{array} \right.$$

The further substitution  $s-1 = e^z$ ,

$$q_0(s, t) = \Psi(z, t) \quad \text{gives}$$

$$\frac{\partial}{\partial t} \Psi(z, t) + k_1 \frac{\partial}{\partial z} \Psi(z, t) = 0$$

whose solution is an arbitrary function of  $(k_1 t - z)$ . For convenience write this as

$$\Psi(z, t) = F \left[ \exp(-k_1 t + z) \right] e^{\frac{-k_2 a}{k_1} z}$$

$$= F \left[ (s-1) e^{-k_1 t} \right] e^{\frac{-k_2 a}{k_1} z} \quad \text{so } \Rightarrow$$

$$\Rightarrow G(s, t) = F \left[ (s-1) e^{-k_1 t} \right] \exp \left[ (s-1) \frac{k_2 a}{k_1} \right]$$

Normalization requires  $G(1, t) = 1$ , and hence

$$F(0) = 1.$$

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The initial conditions determine  $F \Rightarrow$

$$\text{I.C. } p_n(0) = \delta_{n, n_0} \Rightarrow G(s, 0) = F(s-1) e^{(s-1) \frac{k_2 a}{k_1}}$$

$$\Rightarrow G(s, t) = \exp \left[ \frac{k_2}{k_1} a (s-1) (1 - e^{-k_1 t}) \right] \left( 1 + (s-1) e^{-k_1 t} \right)^{n_0}$$

From the generating function we can compute:

$$\langle N(t) \rangle = \partial_s G(s=1, t) = \frac{k_2}{k_1} a (1 - e^{-k_1 t}) + n_0 e^{-k_1 t}$$

$$\langle N(t)^2 \rangle = \frac{\partial_s^2 G(s=1, t)}{\langle N(N-1) \rangle} + \langle N \rangle = \left( n_0 e^{-k_1 t} + \frac{k_2}{k_1} a \right) (1 - e^{-k_1 t})$$

$$\left. \frac{\partial_s^2}{\partial s^2} \right|$$

Moment equations from the differential equation  $\frac{\partial_t G(s, t)}{G(s, t)}$  (\*)

$$k=1, 2, 3, \dots \quad \partial_t \left[ \partial_s^k G(s, t) \right] = \left\{ k \left[ k_2 a \partial_s^{k-1} - k_1 \partial_s^k \right] + (s-1) \left[ k_2 a \partial_s^k - k_1 \partial_s^{k+1} \right] \right\} G(s, t)$$

setting  $s=1$  and using

$$\left. \frac{\partial_s^k}{\partial s^k} \right| \quad \frac{\partial_s^k G(s, t)}{G(s, t)} \Big|_{s=1} = \langle N(t)^k \rangle_f = \langle N(N-1) \dots (N-k+1) \rangle$$

factorial moment

We find

$$\frac{d}{dt} \langle N(t)^k \rangle_f = k \left[ k_2 a \langle N(t)^{k-1} \rangle_f - k_1 \langle N(t)^k \rangle_f \right]$$

and these equations form a closed hierarchy.