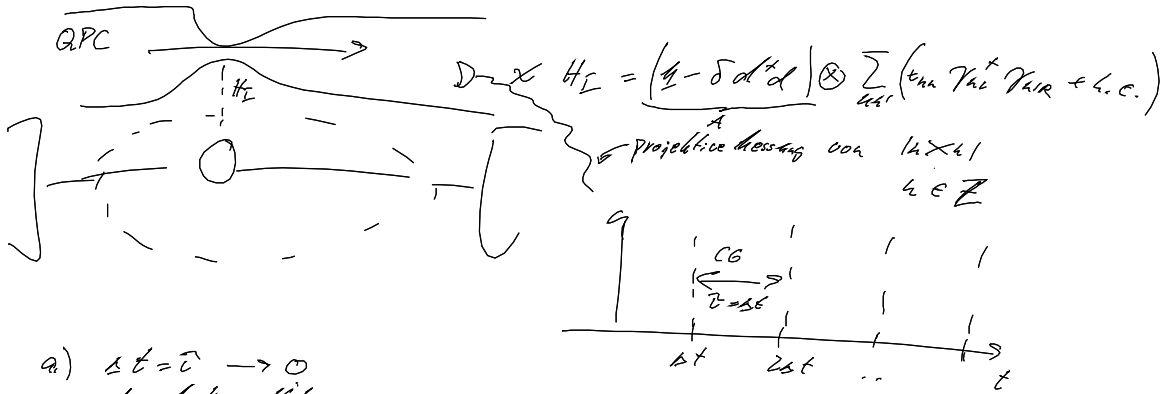


o WdH



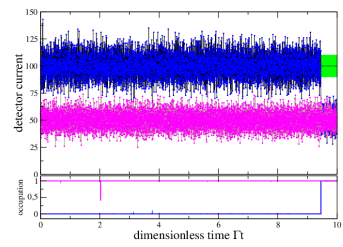
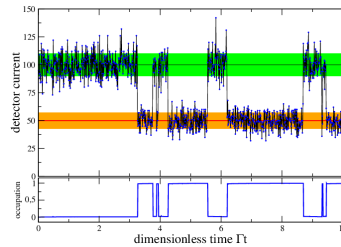
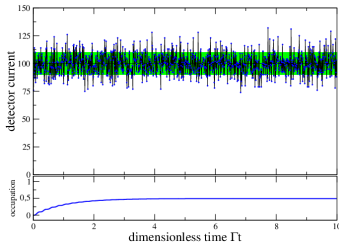
- a) $\Delta t = \bar{c} \rightarrow 0$
- keine Markov-Näherung
 - lokale Lindblad Dynamik

$$\mathcal{L} \rho \sim [e^{ix} A \rho A - \frac{1}{2} \{A^\dagger, \rho\} A]$$

• für $x \rightarrow 0$

$$\mathcal{L} \rho \sim -\gamma [d^\dagger d \rho d d^\dagger + d d^\dagger \rho d^\dagger d]$$

\rightarrow Dämpfung bestimmter Kohärenzen in lok. Basis



b) $\Delta t = \bar{c} \rightarrow \infty$ Markov + Saltz (G) - Näherung

$$\tilde{A}(t) = \hat{A}_0 + \sum_n (\hat{A}_n e^{+i\omega_n t} + \hat{A}_n^\dagger e^{-i\omega_n t})$$

Boson-Frequenzen

$$[\hat{A}_0, \hat{H}_S] = 0$$

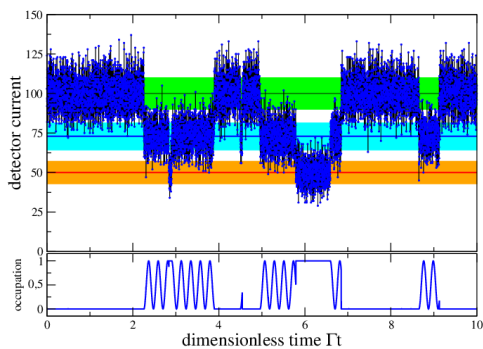
$$\gamma(\omega_n) = 0$$

$$\gamma(0) \neq 0$$

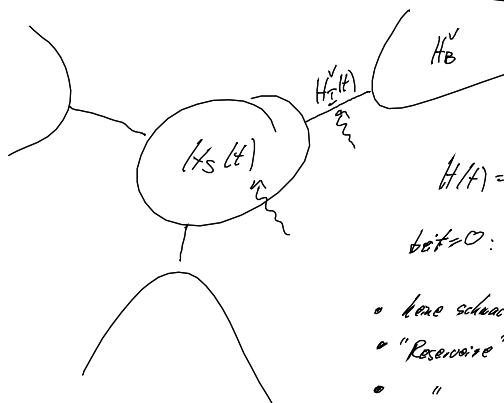
Detektor misst A_0

QND-Messung

→ FT wird durch die Präsenz
eines Det. nicht gestört



7.4. Nicht-perturbative Form der Entropie-Produktion



$$H(A) = H_S(A) + \sum_{\nu} H_{\nu}^{\nu}(A) = \sum_{\nu} H_{\nu}^{\nu}$$

$$\text{bit} \rightarrow 0: \rho_S(0) = \rho_S(0) \otimes \bar{\rho}_{\nu}$$

$$\bar{\rho}_{\nu} = \frac{e^{-\beta_{\nu}(H_{\nu}^{\nu} - \mu_{\nu} N_{\nu}^{\nu})}}{Z_{\nu}}$$

- keine schwache Kopplung
- "Reservoir" können endlich sein
- " " können aus GG getrieben werden

Esposito KJP 2010
van den Broeck
Lindenberg

$$\sum_{\nu} \{ -Tr \{ \rho_{\nu}(t) \ln \rho_{\nu}(t) \} \} = \sum_{\nu} \{ 0 \} = \underbrace{-Tr_S \{ \rho_S(0) \ln \rho_S(0) \}}_{S(0)} - \sum_{\nu} Tr_{\nu} \{ \bar{\rho}_{\nu} \ln \bar{\rho}_{\nu} \}$$

Entropie d. Universums ist konstant

$$\text{red. DA von System } \rho_S(t) = Tr_{\nu} \{ \rho(t) \} \quad \rightarrow \quad S(t) = -Tr_S \{ \rho_S(t) \ln \rho_S(t) \} \text{ nicht konst.}$$

$$\text{Reservoir } \rho_{\nu}(t) = Tr_{S, \nu} \{ \rho(t) \}$$

$$S(0) = \sum_{\nu} \{ -Tr_{\nu} \{ \bar{\rho}_{\nu} \ln \bar{\rho}_{\nu} \} \}$$

$$\begin{aligned} \Delta S(H) &= S(H) - S(0) \\ &= -T_s \{ \rho_s(H) \}_{\text{Gibbs}} + T_r \{ \rho_r(H) \} - \sum_{\nu} T_{\nu} \{ \bar{\rho}_{\nu} \}_{\text{L}} \\ &= -T_r \{ \rho_r(H) \}_{\text{L}} + \dots \\ &= -T_r \{ \rho_r(H) \}_{\text{L}} \left[\rho_s(H) \otimes \bar{\rho}_{\nu} \right] + T_r \{ \rho_r(H) \}_{\text{L}} + \sum_{\nu} T_{\nu} \{ \rho_{\nu}(H) - \bar{\rho}_{\nu} \}_{\text{L}} \\ &= D \left(\rho_r(H) \parallel \rho_s(H) \otimes \bar{\rho}_{\nu} \right) - \sum_{\nu} \beta_{\nu} T_{\nu} \{ \rho_{\nu}(H) - \bar{\rho}_{\nu} \} \left[\langle H_{\nu}^M \rangle - \langle N_{\nu}^M \rangle \right] \\ &= \Delta_i S \geq 0 \qquad \qquad \qquad -\Delta Q^M(H) \end{aligned}$$

$$= \Delta_i S(H) + \sum_{\nu} \beta_{\nu} \cdot \Delta Q^M(H)$$

\uparrow Entropie-Produktion \uparrow Wärme, welche in $[0, t]$ das Reservoir verlässt

$$\Delta_i S(H) = \Delta S(H) - \sum_{\nu} \beta_{\nu} \Delta Q^M(H) \geq 0$$

\uparrow
Änd. der System-Entropie

• $\lim_{t \rightarrow \infty} \frac{1}{t} [\dots]$ • System endlich & keine Stör-ZS an

$$\lim_{t \rightarrow \infty} \frac{\Delta_i S(H)}{t} = - \sum_{\nu} \beta_{\nu} \lim_{t \rightarrow \infty} \frac{\langle (H_{\nu} - \mu_{\nu} N_{\nu}) \rangle_t - \langle (H_{\nu} - \mu_{\nu} N_{\nu}) \rangle_e}{t}$$

Falls gilt

$$\lim_{t \rightarrow \infty} \frac{d}{dt} \langle H_{\nu} \rangle_t = -\frac{T_{\nu}}{E}$$

$$\lim_{t \rightarrow \infty} \frac{d}{dt} \langle N_{\nu} \rangle_t = -\frac{T_{\nu}}{A}$$

$$\Rightarrow \left(- \sum_{\nu} \beta_{\nu} \left(\bar{I}_E^M - \mu_{\nu} \bar{I}_A^M \right) \right) \geq 0$$

im steady state

• für $t > 0$ ist Σ nicht mehr die Summe aus System & Bad-Entropie

$$S_{\nu}(H) = -T_{\nu} \{ \rho_{\nu}(H) \}_{\text{L}} \rho_r(H) \}$$

Korrelations-Entropie

$$S_c(H) \equiv \Sigma(H) - S(H) - \sum_{\nu} S_{\nu}(H)$$

$$-S_c(H) = D \left(\rho_r(H) \parallel \rho_s(H) \otimes \rho_{\nu}(H) \right)$$

man zeigt (Skript) $\Delta_i S(H) \geq -S_c(H) \geq 0$

• $\frac{d}{dt} \Delta_i S(H)$ kann i.A. negativ werden

7.4.1. 2 Qubits

$$H = \frac{h_1}{2} \sigma_1^x + \frac{h_2}{2} \sigma_2^x + \lambda \sigma_1^x \sigma_2^x \quad \sigma_1^x \sigma_2^x = \sigma_1^x \otimes \sigma_2^x$$

$$\rho(0) = \rho_1^0 \otimes \frac{e^{-\beta h_2}}{2} \leftarrow \bar{\rho}_2$$

$$\rho(t) = e^{-iHt} \rho(0) e^{iHt}$$

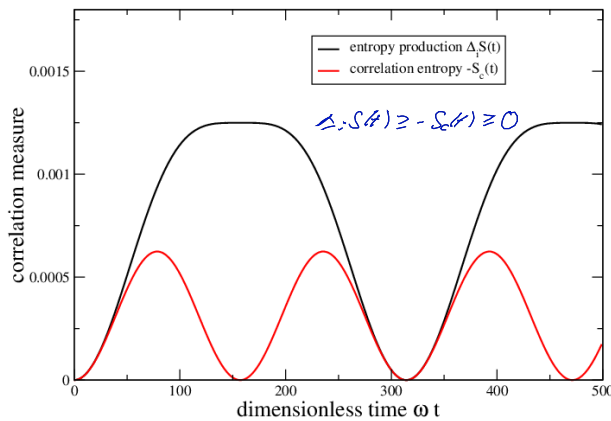
$$\rho_1(t) = \text{Tr}_2 \{ \rho(t) \} \quad \rho_2(t) = \text{Tr}_1 \{ \rho(t) \} = \frac{1}{2} [\mathbb{1} + \vec{h}_2(t) \cdot \vec{\sigma}]$$

$$= \frac{1}{2} [\mathbb{1} + \vec{h}_1(t) \cdot \vec{\sigma}]$$

$$h_1^x = \text{Tr} \{ \sigma_1^x \rho_1(t) \} \quad h_2^x = \text{Tr} \{ \sigma_2^x \rho_2(t) \}$$

$$\Delta S(t) = D(\rho(t) \| \rho_1(t) \otimes \bar{\rho}_2)$$

$$-S_c(t) = D(\rho(t) \| \rho_1(t) \otimes \rho_2(t))$$



7.4.2. Paraphrasing

$$H = \omega \cdot \sigma^z + \lambda \sigma^x \otimes \sum_k (h_k b_k + h_k^* b_k^\dagger) + \sum_k h_k b_k^\dagger b_k$$

$$\rho_0(t) = e^{-iHt} \rho_0^0 \quad f(t) = \frac{t}{\pi} \int_0^\infty \Gamma(\omega) \frac{\sin^2(\omega t/2)}{\omega^2} \text{odd} \left(\frac{h_k}{2} \right) d\omega$$

$$\Delta E(t) = \frac{2}{\pi} \int_0^\infty \frac{\Gamma(\omega)}{\omega} \sin^2(\frac{\omega t}{2}) d\omega$$

$$\Delta H(t) = \frac{2}{\pi} \int_0^\infty \frac{\Gamma(\omega)}{\omega^2} \sin^2(\frac{\omega t}{2}) d\omega$$

} fällt positiv wenn in des Res. abgelesen

$$\Delta S(t) = S(t) - S(0) + \beta [\Delta E(t) - \lambda \Delta H(t)] \geq 0$$

7.4.3. Period. Treiben

$$\Delta S(t) = \Delta_i S(t) + \sum_{\nu} \beta_{\nu} \Delta Q_{\nu}(t)$$

Langzeitwert sei physikalisch periodisch

$$\text{für große } t \quad S(t+T) = S(t)$$

$$\frac{1}{T} \int_t^{t+T} \frac{d}{dt} [\dots]$$

Langzeitstrom seien physikalisch periodisch

$$\text{Periodenmittel} \hat{=} \frac{1}{T} \int_t^{t+T} [\dots] dt'$$

$$0 = \lim_{t \rightarrow \infty} \frac{\Delta_i S(t+T) - \Delta_i S(t)}{T} + \sum_{\nu} \beta_{\nu} \left[\langle I_{E\nu} \rangle - \mu_{\nu} \langle I_{A\nu} \rangle \right]$$

$\xrightarrow{-\langle G \rangle}$

$$\underbrace{\Delta_i S(t+T)}_{\geq 0} = \underbrace{\Delta_i S(t)}_{\geq 0} + \langle G \rangle \cdot T$$

$$\Delta_i S(t+n \cdot T) = \Delta_i S(t) + \langle G \rangle \cdot n \cdot T \quad n \in \mathbb{Z}^+$$

$$\Rightarrow \langle G \rangle = - \sum_{\nu} \beta_{\nu} \left[\langle I_{E\nu} \rangle - \mu_{\nu} \langle I_{A\nu} \rangle \right] \geq 0$$

periodengemittelte Ströme (wenn existent) erfüllen 2. HS