

WdH

$$H_I = \sum_{\alpha} A_{\alpha} \otimes B_{\alpha}$$

System
Bad  
↑
↑  
 $A_{\alpha}^{\dagger}$ 
 $B_{\alpha}^{\dagger}$

• WW-Bild bzgl.  $H_S$  &  $H_B$

$$\tilde{H} = -i \left[ \underbrace{e^{+i H_S + H_B t} H_I e^{-i H_S + H_B t}}_{H_I}, \tilde{P} \right]$$

• Bad-Korrelations-Fkt:

$$C_{\text{op}}(t) = \text{Tr} \left\{ e^{+i H_S t} B_S e^{-i H_B t} B_B \tilde{\rho}_S \right\} \quad \tilde{\rho}_{\text{op}}(t) = \tilde{\rho}_S(t) \otimes \tilde{\rho}_B$$

• Nicht-Markovsche KGL

$$\dot{\tilde{\rho}}_S = - \sum_{\alpha, \beta} \int_0^t dt' \left\{ C_{\text{op}}(t-t') \left[ \tilde{A}_{\alpha}(t), \tilde{A}_{\beta}(t') \right] \tilde{\rho}_S(t') \right\} + \text{h.c.} \quad \text{Tr} \{ \tilde{\rho}_S(t) \} = 1$$

$$= \int_0^t W(t, t') \tilde{\rho}_S(t') dt' \quad \tilde{\rho}_S = \tilde{\rho}_S^{\dagger}$$

↑  
Superoperator

• Markov-Näherung  $C_{\text{op}}(t)$  fallen schnell ab

$$\rightarrow \tilde{\rho}_S(t) \rightarrow \tilde{\rho}_S(t)$$

$$\int_0^t \dots dt' \rightarrow \int_0^{\infty} \dots dt'$$

$$\Rightarrow \dot{\tilde{\rho}}_S = -i [H_S, \tilde{\rho}_S] - \left\{ \sum_{\alpha, \beta} \int_0^{\infty} C_{\text{op}}(\tau) \left[ A_{\alpha}, e^{-i H_B \tau} A_{\beta} e^{+i H_B \tau} \right] \tilde{\rho}_S(t) \right\} + \text{h.c.}$$

↑  
Korrelations im Lindblad-Bild

$$= W \tilde{\rho}_S(t) \quad \rightarrow \quad \tilde{\rho}_S(t) = e^{W t} \tilde{\rho}_S^0$$

$$\neq \text{Tr} \{ \tilde{\rho}_S \} = 1 \quad \text{h.c.} \quad \tilde{\rho}_S = \tilde{\rho}_S^{\dagger} \quad \neq \mathcal{E}$$

aber keine Lindblad-Form (in. Allg.)

$$B_S \rho_B = \frac{c_1}{2} b_1^{\dagger} + \frac{c_2}{2} b_2^{\dagger} + 2 \sqrt{c_1 c_2} b_1^{\dagger} b_2^{\dagger}$$

a.) Säkular-Näherung

$$\dot{\tilde{\rho}}_S = \sum_{\omega} \sum_{\alpha} Z_{\omega \alpha} e^{i(\omega - \omega_{\alpha})t} \tilde{\rho}_S(t) = \sum_{\alpha} Z_{\omega \alpha} \tilde{\rho}_S(t) + \sum_{\omega \neq \omega_{\alpha}} Z_{\omega \alpha} e^{+i(\omega - \omega_{\alpha})t} \tilde{\rho}_S(t)$$

↑
↑  
Super-Operatoren
Eigen-Frequenzen des Systems (Kügelige)

( "Säkular-Näherung" )

$$\dot{\tilde{\rho}}_S = - \sum_{\alpha, \beta} \int_0^{\infty} \left\{ C_{\text{op}}(\tau) \left[ \tilde{A}_{\alpha}(t), \tilde{A}_{\beta}(t-\tau) \right] \tilde{\rho}_S(t) \right\} + \text{h.c.}$$

$$e^{+i k_1 t} A_{\alpha} e^{-i k_2 t} = \sum_{\omega, \delta} e^{+i E_{\alpha} t} |a \times a| A_{\alpha} |b \times b| e^{-i E_{\beta} t}$$

$$H_S = \sum_{\alpha} E_{\alpha} |a \times a| = \sum_{\omega, \delta} e^{+i(E_{\alpha} - E_{\beta})t} \langle a | A_{\alpha} | b \rangle |a \times b|$$

→ Gleichung mit  $\delta_{E_{\alpha} - E_{\beta}, E_{\alpha} - E_{\beta}}$  führt zwar auf Lindblad-Form

$$\text{Halbspektrale FT: } \Gamma_{\text{op}}(\omega) = \int_0^{\infty} C_{\text{op}}(\tau) e^{+i \omega \tau} d\tau = \frac{1}{2} \mathcal{F}_{\text{op}}(\omega) + \frac{1}{2} \mathcal{F}_{\text{op}}^{\dagger}(\omega)$$

$$\mathcal{F}_{\text{op}}(\omega) = \Gamma_{\text{op}}(\omega) + \Gamma_{\text{op}}^{\dagger}(\omega) = \int_{-\infty}^{\infty} C_{\text{op}}(\tau) e^{+i \omega \tau} d\tau \quad \text{volle FT (gerade)}$$

$$\mathcal{F}_{\text{op}}^{\dagger}(\omega) = \Gamma_{\text{op}}(\omega) - \Gamma_{\text{op}}^{\dagger}(\omega) = \int_{-\infty}^{\infty} C_{\text{op}}(\tau) \cdot \text{sgn}(\tau) e^{+i \omega \tau} d\tau \quad \text{volle ungerade FT}$$

$$= \frac{i}{\pi} \oint \frac{\gamma_{rs}(\omega')}{\omega - \omega'} d\omega'$$

Cauchy Hauptwert

Vorr:  $H_S = \sum A_r \otimes B_r$      $A_r = A_r^\dagger$      $B_r = B_r^\dagger$

$[\hat{H}_S, \hat{\rho}_S] = 0$      $\text{Tr}\{B_r \hat{\rho}_S\} = 0$      $C_{rs}(\omega) = \text{Tr}\{e^{+i\omega \hat{H}_S} B_r e^{-i\omega \hat{H}_S} B_s - B_s\}$

berechne  $\gamma_{rs}(\omega) = \int C_{rs}(\omega) e^{+i\omega \hat{H}_S} d\omega$      $\hat{H}_S(\omega)$  analog

$\dot{\hat{\rho}}_S = -i[\hat{H}_S + \sum_{ab} \hat{G}_{ab} \hat{L}_{ab}, \hat{\rho}_S] + \sum_{ab,cd} \gamma_{ab,cd} [\hat{L}_{ab} \hat{\rho}_S \hat{L}_{cd}^\dagger - \frac{1}{2} \{\hat{L}_{ab}^\dagger \hat{L}_{ab}, \hat{\rho}_S\}]$

$\hat{L}_{ab} = |a\rangle \langle b|$      $H_S |a\rangle = E_a |a\rangle$

$\hat{G}_{ab} = \sum_{rs} \sum_0 \frac{1}{2i} \hat{G}_{rs} (E_b - E_a) \delta_{E_b - E_a, E_r - E_s} \langle c|A_r|b\rangle \langle c|A_s|a\rangle^*$

$\gamma_{ab,cd} = \sum_{rs} \gamma_{rs} (E_b - E_a) \delta_{E_b - E_a, E_r - E_s} \langle a|A_r|b\rangle \langle c|A_s|d\rangle^*$

BHS Hermitengleichung

+ zeige:  $\hat{H}_S = \sum_{ab} \hat{G}_{ab} \hat{L}_{ab} = \hat{H}_S^\dagger$

$-\sum_{ab,cd} x_{ab}^* \gamma_{ab,cd} x_{cd} \geq 0$

folgt aus  $\gamma_{rs}(\omega)$  ist pos. definit

- +  $\hat{L}_{ab}$  sind nicht lokal
- + oft wird  $H_S$  vernachlässigt
- + falls  $\hat{\rho}_S$   $N \times N \rightarrow \mathcal{L} \hat{\rho}_S = \dot{\hat{\rho}}_S$      $\mathcal{L} : N^2 \times N^2$
- + falls  $H_S$  keine Entartungen hat

$\rightarrow \delta_{E_a, E_b} = \delta_{ab}$

$\rightarrow \dot{\hat{\rho}}_{aa} = + \sum_b \gamma_{ab,ab} \hat{\rho}_{bb} - \left( \sum_b \gamma_{ba,ba} \right) \hat{\rho}_{aa}$     Rotationsgleichung

$\gamma_{ab,ab} = \sum_{rs} \gamma_{rs} (E_b - E_a) \langle a|A_r|b\rangle \langle a|A_s|b\rangle^* \geq 0$

pos. definit

$\mathcal{L}_{BAS} = \begin{pmatrix} \mathcal{L}_{BAS} & | & \textcircled{+} \\ \hline \textcircled{+} & | & \mathcal{L}_{cot} \end{pmatrix}$      $\hat{\rho} = \begin{pmatrix} \hat{\rho}_{BAS} \\ \hline -\hat{\rho}_{cot} \\ \vdots \end{pmatrix}$

+ Lindblad Davis - Abb:  $\mathcal{L}_{BAS} e^{-B \hat{H}_S} = 0$

für  $\hat{\rho}_B = \frac{e^{-\beta \hat{H}_S}}{Z_B}$

1.3.3. Harmon. Oszillator in Stam. Bad

$H_{tot} = \underbrace{\alpha a^\dagger a}_{H_S} + \underbrace{\sum_k \omega_k b_k^\dagger b_k}_{H_B} + \underbrace{(a + \epsilon a^\dagger) \sum_k \frac{1}{k} (h_k b_k + h_k^\dagger b_k^\dagger)}_{H_C}$      $A = A^\dagger$   
 $B = B^\dagger$

$\dot{\hat{\rho}} = -i[\hat{H}_T, \hat{\rho}]$      $\hat{H}_T = (a \cdot e^{-i\omega t} + a^\dagger \cdot e^{+i\omega t}) \otimes \sum_k (h_k b_k e^{-i\omega_k t} + h_k^\dagger b_k^\dagger e^{+i\omega_k t})$

$\tilde{a}(t) = e^{+i\omega t a^\dagger} a e^{-i\omega t a}$

$\rightarrow \frac{d}{dt} \tilde{a}(t) = -i\omega \tilde{a}(t)$

$\alpha = \beta = 1$

$$\begin{aligned}
 \text{Tr}\{B \bar{\rho}_B\} &= 1 \\
 C_{\mu\nu}(\omega) &= C(\omega) = \text{Tr} \left\{ \frac{1}{2} (b_{\mu} b_{\nu} e^{-i\omega\bar{v}} + b_{\nu}^{\dagger} b_{\mu}^{\dagger} e^{+i\omega\bar{v}}) \left( \sum_{\vec{k}} b_{\vec{k}} b_{\vec{k}}^{\dagger} + b_{\vec{k}}^{\dagger} b_{\vec{k}} \right) \frac{e^{-\beta \sum_{\vec{k}} \omega_{\vec{k}} b_{\vec{k}}^{\dagger} b_{\vec{k}}}}{Z_B} \right\} \\
 &= \frac{1}{2\pi} \int_0^{\infty} f(\omega) \left[ e^{-i\omega\bar{v}} [1 + b_{\vec{k}}(\omega)] + e^{+i\omega\bar{v}} \cdot b_{\vec{k}}(\omega) \right] d\omega \\
 \tilde{f}(\omega) &= -\tilde{f}(-\omega) & f(\omega) &= \frac{1}{2\pi} \sum_{\vec{k}} \omega_{\vec{k}}^2 \cdot \delta(\omega - \omega_{\vec{k}}) \\
 \tilde{f}(\omega) &= f(\omega) & & \text{Spektrale Dichte} \\
 &= \frac{1}{2\pi} \int \underbrace{\tilde{f}(\omega) [1 + b_{\vec{k}}(\omega)]}_{\gamma_{\vec{k}}(\omega)} e^{-i\omega\bar{v}} d\omega
 \end{aligned}$$

$$\dot{\rho} = -i [\rho, a^{\dagger} a] - \int C(\bar{v}) [a + a^{\dagger}, e^{-i\omega a^{\dagger} a \bar{v}} (a + a^{\dagger}) e^{+i\omega a^{\dagger} a \bar{v}} \rho(t)] d\bar{v} + \text{h.c.}$$

Karboon-ME (keine Lindblad-Form)

vernachlässige  $e^{\pm 2i\omega t}$  im MW-Bild

$$\begin{aligned}
 \Gamma(\omega) &= \int_0^{\infty} C(\bar{v}) e^{+i\omega\bar{v}} d\bar{v} \Rightarrow \frac{1}{2} \gamma + \frac{i}{2} \beta = \Gamma(+\omega) \\
 \frac{1}{2} \tilde{\gamma} + \frac{i}{2} \tilde{\beta} &= \Gamma(-\omega)
 \end{aligned}$$

$$\Rightarrow \dot{\rho} = -i \left[ \frac{\beta}{2} a^{\dagger} a + \frac{\tilde{\beta}}{2} a a^{\dagger}, \rho \right] + \gamma [a^{\dagger} a \rho a^{\dagger} - \frac{1}{2} \{a^{\dagger} a, \rho\}] + \tilde{\gamma} [a^{\dagger} \rho a - \frac{1}{2} \{a a^{\dagger}, \rho\}]$$

$$\Rightarrow \dot{\rho} = -i \left[ \omega a^{\dagger} a + \left( \frac{\beta}{2} + \frac{\tilde{\beta}}{2} \right) a^{\dagger} a, \rho \right] + \gamma [a^{\dagger} a \rho a^{\dagger} - \frac{1}{2} \{a^{\dagger} a, \rho\}] + \tilde{\gamma} [a^{\dagger} \rho a - \frac{1}{2} \{a a^{\dagger}, \rho\}]$$

$$H_0 |n\rangle = \omega n |n\rangle = E_n |n\rangle$$

$$a^{\dagger} a = \sum_{k=1}^{\infty} k \cdot |k\rangle \langle k|$$

$$a = \sum_{k=1}^{\infty} \sqrt{k} |k-1\rangle \langle k|$$

### 1.3.4. Gleichgewichts-TD

$$\text{Wann ist } \sum e^{-\beta k \omega} = 0$$

$$\text{falls } \rho_0 = \frac{e^{-\beta H_0}}{Z_B}$$

$$C_{\beta, \mu\nu}(\bar{v}) = C_{\beta, \nu\mu}(-\bar{v} - i\beta) \quad \text{Kubo-Martin-Schwinger Relation (KMS)}$$

$$\begin{aligned}
 C_{\beta, \mu\nu}(-\bar{v} - i\beta) &= \text{Tr} \left\{ \frac{e^{+i\omega_0(-\bar{v} - i\beta)}}{B_{\beta}} B_{\beta} e^{-i\omega_0(-\bar{v} - i\beta)} B_{\beta} \frac{e^{-\beta H_0}}{Z_B} \right\} \\
 &= \text{Tr} \left\{ \frac{e^{-\beta H_0}}{Z_B} \frac{e^{+i\omega_0 \bar{v}}}{B_{\beta}} B_{\beta} \frac{e^{-\beta H_0}}{Z_B} \frac{e^{+i\omega_0 \bar{v}}}{B_{\beta}} \frac{e^{-i\omega_0 \bar{v}}}{B_{\beta}} B_{\beta} \right\} = \text{Tr} \left\{ B_{\beta}(\bar{v}) B_{\beta} \frac{e^{-\beta H_0}}{Z_B} \right\} \\
 &= C_{\beta, \nu\mu}(\bar{v})
 \end{aligned}$$

$$\frac{k_B(\omega)}{1 + k_B(\omega)} = e^{-\beta \omega}$$

Wegen: 2. Hauptsatz