

WdK
WZ-Verteilung

$$W^i(t) = \frac{\text{Tr} \{ Z_i \cdot e^{z_i \cdot \hat{O}} \cdot Z_i^\dagger \}}{\text{Tr} \{ Z_i \cdot \hat{\rho} \}}$$

\uparrow 2. Sprung \uparrow 1. Sprung
 \uparrow kond. auf 1. Sprung \uparrow langzeit-WZ-Verteilung

mittler WZ zu Sprüngen (j) and (i)

$$k\text{-tes Moment } \langle \hat{O}^k \rangle = \int_0^\infty W^i(t) \cdot \hat{O}^k dt$$

• Mikroskop Ableitung von ZF 2-Messungen Bed: $t_0=0$
 $t_1=t$



Messe der Differenz einer Bed-Observable \hat{O} Wert der 1. Messung

$$M_{\text{stat}}(x,t) = \sum_z \text{Tr} \{ e^{-z \hat{O}} (\hat{O} - O_0) \rho_S(t) \otimes \hat{\rho}_B^{(1)} M_{\text{stat}}^+(t) \}$$

Mittlung über auf Messgeschichte

$$12 \times 12 \hat{\rho}_B \otimes 12 \times 12 = \hat{\rho}_B^{(1)}$$

kond. auf Mess. 1

$$\text{z.B. } (-i \hat{O}_z)^k M_{\text{stat}}(z,t) \Big|_{z=0} = \sum_P \text{Tr} \{ (\hat{O} - O_0)^k \rho_{\text{stat}}(t) \}$$

→ verallg. Zeit-Evol.-Op. $M_{\text{stat}}(t) = e^{-i \hat{O} z t} U(t) e^{-i \hat{O} z t} \sum_P \hat{\rho}_B^{(1)}$

→ Abl. der Messgl. z.B.: $[M_z + t \cdot z + (z) \dots] \rho_S \stackrel{!}{=} \text{Tr}_B \{ M_{\text{stat}}(t) \rho_S \otimes \hat{\rho}_B M_{\text{stat}}^+(t) \}$

→ in der Messgleichung bekommt die Stringen ein Zeitfeld "L_z ∈ L_B⁺ → L_z ∈ L_B⁺ · e^{+iz}"

→ verallg. Korrelationsfunktion

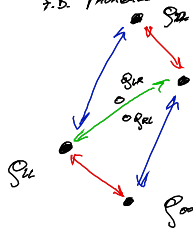
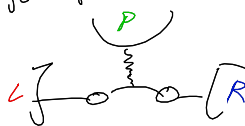
$$C_{\text{op}}(t) \rightarrow C_{\text{op}}^z(t) = \text{Tr}_B \{ e^{-i \hat{O} x} e^{+i t_0 \hat{P}} \rho_B e^{-i t_0 \hat{O}} e^{+i \hat{O} x} \rho_B \cdot \hat{\rho}_B \}$$

$$\text{falls } \hat{\rho}_B = \sum_P P_e \text{ } 12 \times 12 \text{ } \Leftrightarrow [\hat{\rho}_B, \hat{O}] = 0$$

$$\rightarrow \hat{\rho}_B = \hat{\rho}_B$$

Wann sinnvoll

+ Kohärenzen
z.B. Phänomen-assul. Tunnel



Treading mit mikroskop. ZF möglich

3.4.1. Beispiel: SRL $\hat{O} = H_0$

$$H = \underbrace{\varepsilon d^\dagger d}_{H_S} + d \otimes \sum_k \underbrace{t_k}_{A_1} c_k^\dagger + d \otimes \sum_k \underbrace{t_k^*}_{A_2} c_k + \sum_k \underbrace{\varepsilon_k}_{B_0} c_k^\dagger c_k$$

$$\begin{aligned} e^{+i\alpha c_k^\dagger c_k} c_k^\dagger e^{-i\alpha c_k^\dagger c_k} \\ = c_k^\dagger \cdot e^{+i\alpha} \end{aligned}$$

$$\begin{aligned} C_{12}^z(\tau) &= \sum_{k_1, k_2} \text{Tr}_B \left\{ t_{k_1} c_{k_1}^\dagger e^{+i\varepsilon_{k_1}\tau} e^{-i\varepsilon_{k_2}\tau} t_{k_2}^* c_{k_2} \bar{\rho}_B \right\} \\ &= \sum_k (t_k)^2 \cdot f(\varepsilon_k) \cdot e^{+i\varepsilon_k(\tau-\tau)} \\ &= \frac{1}{2\pi} \int \Gamma(\omega) f(\omega) e^{+i\omega\tau} e^{-i\omega\tau} d\omega \\ &= \frac{1}{2\pi} \int \Gamma(\omega) \cdot f(\omega) e^{+i\omega\tau} e^{-i\omega\tau} d\omega \\ &\quad \Gamma_{12}^z(\omega) \end{aligned}$$

analog: $C_{21}^z(\tau) = \frac{1}{2\pi} \int \Gamma(\omega) [1-f(\omega)] e^{+i\omega\tau} e^{-i\omega\tau} d\omega$
 $\Gamma_{21}^z(\omega)$

→ BHS Perturbations (oder GG mit $\tau \rightarrow \infty$)

$$\Gamma_{ab,ab}^z = \sum_{\gamma, \beta} \Gamma_{\gamma\beta}^z (E_\beta - E_\alpha) \langle a | A_\alpha | b \rangle \langle a | A^\dagger | b \rangle$$

Rate $b \rightarrow a$

$$E_0 = 0 \quad E_1 = \varepsilon$$

$$\begin{aligned} \Gamma_{01,01}^z &= \Gamma_{11}^z(+\varepsilon) = \Gamma(\varepsilon) [1-f(\varepsilon)] e^{+\varepsilon\tau} \\ \Gamma_{10,10}^z &= \Gamma_{11}^z(-\varepsilon) = \Gamma(\varepsilon) f(\varepsilon) \cdot e^{-i\varepsilon\tau} \end{aligned}$$

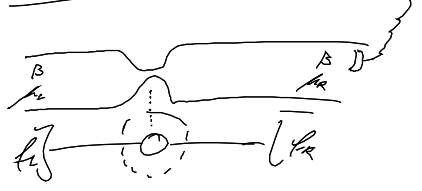
Konvention: positiv, wenn in die Rec gerichtet

$$\hookrightarrow Z(\varepsilon) = \Gamma(\varepsilon) \begin{pmatrix} -f(\varepsilon) & +[1-f(\varepsilon)] e^{+\varepsilon\tau} \\ f(\varepsilon) \cdot e^{-i\varepsilon\tau} & -[1-f(\varepsilon)] \end{pmatrix}$$

→ konsistent

analog $\hat{O} = H_0$

3.4.2. Beispiel QPC $\hat{O} = N_{B,QPC}$



$$H_{SET} = \varepsilon d^\dagger d + \sum_{\substack{a,v \\ \text{verteilt}}} (t_{av} d c_{av}^\dagger + t_{av}^* c_{av}^\dagger d) + \sum_{\alpha, \nu} \varepsilon_{\alpha\nu} c_{\alpha\nu}^\dagger c_{\alpha\nu}$$

$$H_{QPC} = \sum_k \sum_{\nu, \alpha} \varepsilon_{k\nu} \gamma_{k\nu}^\dagger \gamma_{k\nu}$$

$$H_{\text{SET-QPC}} = \underbrace{\left(1 - i \cdot \text{d.t.d.}\right)}_A \sum_{kk'} \underbrace{\left(t_{kk'} \gamma_{kL} \gamma_{kR}^+ + t_{kk'}^* \gamma_{kR} \gamma_{kL}^+ \right)}_B$$

Reduktion des Terms bei besetzten SET-QD

$$0 < \nu < 2$$

betrachte nur die Kopplung an das QPC

$$\hat{O} = \sum_n \gamma_{nL}^+ \gamma_{nR}$$

$$\rho_{B,QPC} = \frac{e^{-\beta \sum_k (E_{kL} - \mu_L) \gamma_{kL}^+ \gamma_{kL}}}{Z_L} \frac{e^{-\beta \sum_k (E_{kR} - \mu_R) \gamma_{kR}^+ \gamma_{kR}}}{Z_R}$$

$$C^q(\omega) = \text{Tr} \left\{ e^{-i\eta \hat{O}} e^{+i\eta \text{trac} \hat{O}} B e^{-i\eta \text{trac} \hat{O}} e^{+i\eta \hat{O}} B \bar{\rho}_B \right\} \quad (\bar{\rho}_B = \bar{\rho}_0)$$

$$= \sum_{kk'} \text{Tr} \left\{ \left[t_{kk'} \gamma_{kL} \gamma_{kR}^+ e^{+i(E_{kR} - E_{kL})\nu} e^{-i\eta} + t_{kk'}^* \gamma_{kR} \gamma_{kL}^+ e^{-i(E_{kR} - E_{kL})\nu} e^{+i\eta} \right] \right.$$

$$\left. \times \left[t_{kk'} \gamma_{kL} \gamma_{kR}^+ + t_{kk'}^* \gamma_{kR} \gamma_{kL}^+ \right] \rho_{B,QPC} \right\}$$

$$= \sum_{kk'} \left\{ e^{-i\eta} |t_{kk'}|^2 e^{+i(E_{kR} - E_{kL})\nu} [1 - f_L(E_{kL})] \cdot f_R(E_{kR}) \right.$$

$$\left. + e^{+i\eta} |t_{kk'}|^2 e^{-i(E_{kR} - E_{kL})\nu} f_L(E_{kL}) [1 - f_R(E_{kR})] \right\}$$

$$T(\omega, \omega') = 2\pi \sum_{kk'} |t_{kk'}|^2 \delta(\omega - E_{kL}) \delta(\omega' - E_{kR}) \quad \text{Transmission}$$

$$\gamma(\omega) = \int C(\omega) e^{+i\omega \nu} d\omega$$

$$\rightarrow = e^{-i\eta} \frac{1}{2\pi} \int d\omega \int d\omega' T(\omega, \omega') e^{+i(\omega - \omega')\nu} [1 - f_L(\omega)] f_R(\omega')$$

$$+ e^{+i\eta} \frac{1}{2\pi} \int d\omega \int d\omega' T(\omega, \omega') e^{-i(\omega - \omega')\nu} f_L(\omega) [1 - f_R(\omega')] \quad \frac{1}{2\pi} \int e^{i\omega \nu} d\omega \delta(\omega)$$

$$\gamma^q(\omega) = e^{-i\eta} \int d\omega T(\omega, \omega - \omega) [1 - f_L(\omega)] f_R(\omega - \omega) \quad \text{Transfer von } R \rightarrow L$$

mit Absorption von Energie ω aus der Spalte

$$+ e^{+i\eta} \int d\omega T(\omega, \omega + \omega) f_L(\omega) [1 - f_R(\omega + \omega)] \quad \text{Transfer von } L \rightarrow R$$

Integrale des Typs $I = \int d\omega f_L(\omega) [1 - f_R(\omega)]$

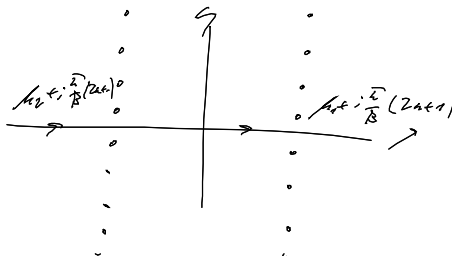
Vereinfachung $T(\omega, \omega') \approx T$

a) für $\beta \rightarrow \infty$ (Temp. $\rightarrow 0$)

$$I \rightarrow (f_L - f_R) \cdot \Theta(\mu_L - \mu_R)$$

b) für endl. Temp.

Funktionentheorie



$$\dots T = \frac{A_1 A_2}{1 - e^{-\beta(E_1 - E_2)}}$$

$$\gamma_{ab}^z = \gamma^z (E_a - E_b) |\langle a | [\hat{L} - \tilde{U} a^\dagger a] | b \rangle|^2$$

$$\Sigma_{QPC}(\xi) = \begin{pmatrix} 1 & 0 \\ 0 & (1-\alpha)^2 \end{pmatrix} T \left[(e^{-\beta\xi} - 1) \frac{V}{e^{\beta V} - 1} + (e^{\beta\xi} - 1) \frac{V}{1 - e^{-\beta V}} \right]$$

$$V = \mu_L^{QPC} - \mu_R^{QPC}$$

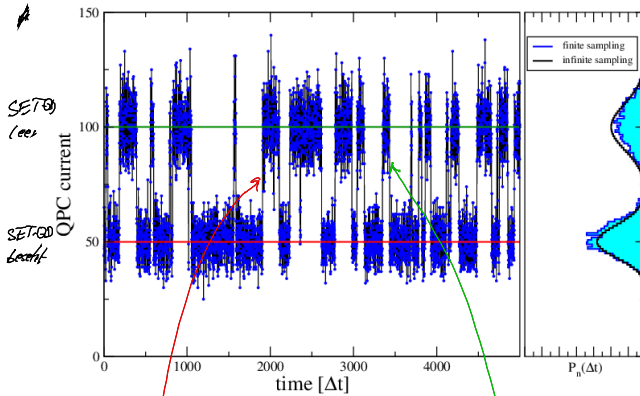
$$\Sigma_{ges}(\xi) = \Sigma_{SET} + \Sigma_{QPC}(\xi)$$

a.) z.B. $\Gamma_L = \Gamma_R \rightarrow 0$

$$\begin{matrix} \uparrow & \uparrow \\ \bar{I}_0 = T \cdot V & \bar{I}_1 = (1-\alpha)^2 \cdot \bar{I}_0 \end{matrix}$$

SET-QD leer falls: $\{\Gamma_L, \Gamma_R\} \ll \{T \cdot V, (1-\alpha)^2 \cdot T \cdot V\}$

b) $P_{kur}(t) = \frac{1}{2\pi} \int Tr \{ e^{i \Sigma_{ges}(\xi) \cdot t - i \xi \xi} \} d\xi$



$$\bar{I}_{QPC} = \frac{\Sigma R}{\Delta \phi}$$

SET entleert sich (ins rechte Bad)

SET füllt sich (aus dem linken Bad)