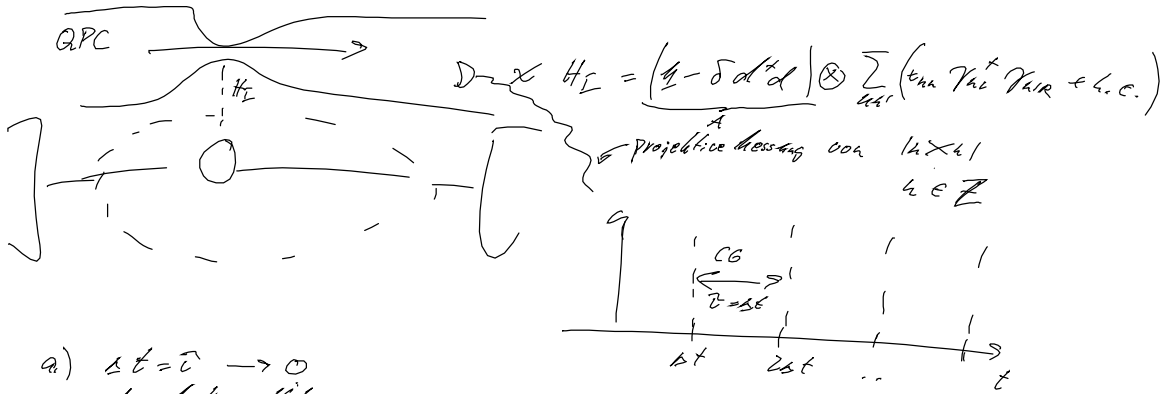


o WdH



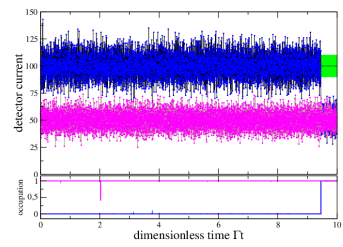
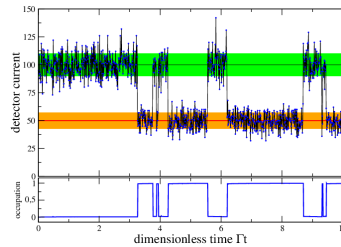
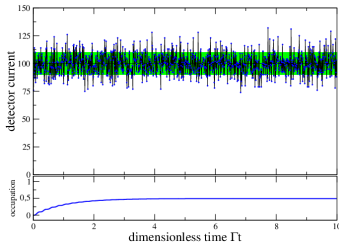
- a) $\Delta t = \bar{c} \rightarrow 0$
- keine Markov-Näherung
 - lokale Lindblad Dynamik

$$\mathcal{L} \rho \sim [e^{ix} A \rho A - \frac{1}{2} \{A^\dagger, \rho\} A]$$

• für $x \rightarrow 0$

$$\mathcal{L} \rho \sim -\gamma [d^\dagger d \rho d d^\dagger + d d^\dagger \rho d^\dagger d]$$

\rightarrow Dämpfung bestimmter Kohärenzen in lok. Basis



b) $\Delta t = \bar{c} \rightarrow \infty$ Markov + Sättigung-Näherung

$$\tilde{A}(t) = \hat{A}_0 + \sum_n (\hat{A}_n e^{+i\omega_n t} + \hat{A}_n^\dagger e^{-i\omega_n t})$$

Boson-Frequenzen

$$[\hat{A}_0, \hat{H}_S] = 0$$

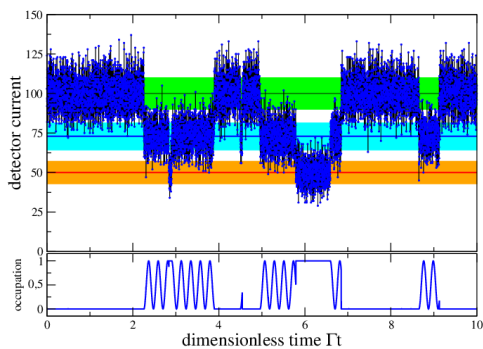
$$\gamma(\omega_n) = 0$$

$$\gamma(0) \neq 0$$

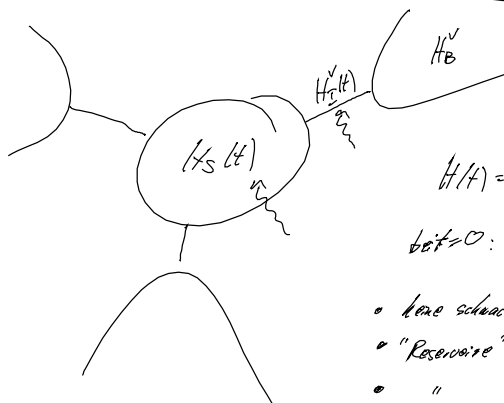
Defektor misst A_0

QND-Messung

→ FT wird durch die Präsenz
eines Det. nicht gestört



7.4. Nicht-perturbative Form der Entropie-Produktion



$$H(t) = H_S(t) + \sum_{\nu} H_{\nu}^V(t) + \sum_{\nu} H_B^V$$

$$\text{bit} \neq 0: \rho(t) = \rho_S(t) \otimes \bar{\rho}_V$$

$$\bar{\rho}_V = \frac{e^{-\beta_V(H_B^V - \mu_V N_B^V)}}{Z_V}$$

- keine schwache Kopplung
- "Reservoir" können endlich sein
- " " können aus GG getrieben werden

Esposito KJP 2010
van den Broeck
Lindenberg

$$\sum_{\nu} \{ -\text{Tr} \{ \rho_{\nu}(t) \ln \rho_{\nu}(t) \} \} = \sum_{\nu} \{ 0 \} = \underbrace{-\text{Tr}_S \{ \rho_S(t) \ln \rho_S(t) \}}_{S(t)} - \sum_{\nu} \text{Tr}_V \{ \bar{\rho}_V \ln \bar{\rho}_V \}$$

Entropie d. Universums ist konstant

$$\text{red. DA von System } \rho_S(t) = \text{Tr}_V \{ \rho(t) \} \quad \rightarrow \quad S(t) = -\text{Tr}_S \{ \rho_S(t) \ln \rho_S(t) \} \text{ nicht konst.}$$

$$\text{Reservoir } \rho_V(t) = \text{Tr}_S \{ \rho(t) \}$$

$$S(t) = \sum_{\nu} \{ -\text{Tr}_V \{ \bar{\rho}_V \ln \bar{\rho}_V \} \}$$

$$\begin{aligned}
\Delta S(H) &= S(H) - S(0) \\
&= -T_S \{ \rho_S(H) \} \underbrace{\ln \rho_S(H)}_{\text{ent. ins}} + T_V \{ \rho_V(H) \} \ln \rho_V(H) - \sum_V T_V \{ \bar{p}_V \} \ln \bar{p}_V \\
&= -T_V \{ \rho_V(H) \} \ln \rho_V(H) + \dots \\
&= -T_V \{ \rho_V(H) \} \ln \left[\rho_V(H) \otimes \bar{p}_V \right] + T_V \{ \rho_V(H) \} \ln \rho_V(H) + \sum_V T_V \{ \bar{p}_V \} \ln \bar{p}_V \\
&= \underbrace{D(\rho_V(H) \parallel \rho_V(H) \otimes \bar{p}_V)}_{\Delta S} - \sum_V \beta_V T_V \{ \bar{p}_V \} \underbrace{[H_B^{(M)} - \mu_B^{(M)}]}_{-\Delta Q^{(M)}(H)}
\end{aligned}$$

$$\Delta S(H) = \Delta S(H) + \sum_V \beta_V \Delta Q^{(M)}(t)$$

↑ Entropie-Produktion ↑ Wärme, welche in $[0, t]$ das Reservoir verlässt

$$\Delta S(H) = \Delta S(H) - \sum_V \beta_V \Delta Q^{(M)}(H) \geq 0$$

↑
Änd. der System-Entropie

• $\lim_{t \rightarrow \infty} \frac{1}{t} [\dots]$ • System endlich & keine Stör-ZS an

$$\lim_{t \rightarrow \infty} \frac{\Delta S(H)}{t} = - \sum_V \beta_V \lim_{t \rightarrow \infty} \frac{\langle (H_V - \mu_V N_V) \rangle_t - \langle (H_V - \mu_V N_V) \rangle_0}{t}$$

$$\Rightarrow \left(- \sum_V \beta_V (\bar{I}_E^{(M)} - \mu_V \bar{I}_A^{(M)}) \right) \geq 0$$

im steady state

Falls gilt

$$\lim_{t \rightarrow \infty} \frac{d}{dt} \langle H_V \rangle_t = -\frac{1}{\tau} \langle H_V \rangle_t$$

$$\lim_{t \rightarrow \infty} \frac{d}{dt} \langle N_V \rangle_t = -\frac{1}{\tau} \langle N_V \rangle_t$$

• für $t > 0$ ist Σ nicht mehr die Summe aus System & Bad-Entropie

$$S_V(H) = -T_V \{ \rho_V(H) \} \ln \rho_V(H)$$

Korrelations-Entropie

$$S_C(H) \equiv \Sigma(H) - S(H) - \sum_V S_V(H)$$

$$-S_C(H) = D(\rho(H) \parallel \rho_S(H) \otimes \rho_V(H))$$

man zeigt (Skript) $\Delta S(H) \geq -S_C(H) \geq 0$

• $\frac{d}{dt} \Delta S(H)$ kann i. A. negativ werden

7.4.1. 2 Qubits

$$H = \frac{h_1}{2} \sigma_1^z + \frac{h_2}{2} \sigma_2^z + \lambda \sigma_1^x \sigma_2^x \quad \sigma_1^x \sigma_2^x = \sigma_1^x \otimes \sigma_2^x$$

$$\rho(0) = \rho_1^0 \otimes \frac{e^{-\beta h_2}}{Z_2} \leftarrow \bar{\rho}_2$$

$$\rho(t) = e^{-iHt} \rho(0) e^{iHt}$$

$$\rho_1(t) = \text{Tr}_2 \{ \rho(t) \}$$

$$= \frac{1}{2} [I_2 + \vec{r}_1(t) \cdot \vec{\sigma}]$$

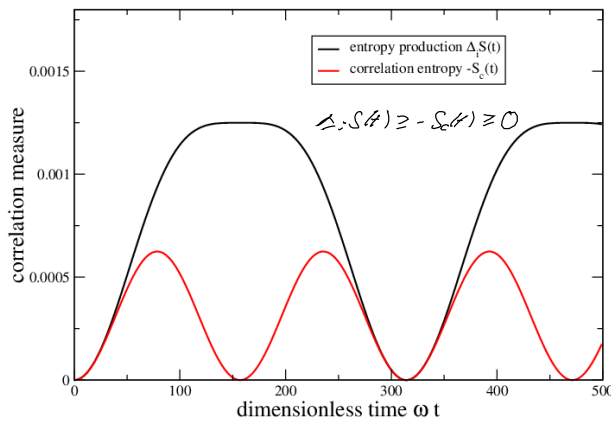
$$\rho_2(t) = \text{Tr}_1 \{ \rho(t) \} = \frac{1}{2} [I_2 + \vec{r}_2(t) \cdot \vec{\sigma}]$$

$$h_1^x = \text{Tr} \{ \sigma_1^x \rho_1(t) \}$$

$$h_2^x = \text{Tr} \{ \sigma_2^x \rho_2(t) \}$$

$$\Delta S(t) = D(\rho(t) \| \rho_1(t) \otimes \bar{\rho}_2)$$

$$-S_c(t) = D(\rho(t) \| \rho_1(t) \otimes \rho_2(t))$$



7.4.2. Pure dephasing

$$H = \omega_L \sigma^z + \lambda \sigma^z \otimes \sum_k (\lambda_k b_k + \lambda_k^* b_k^\dagger) + \sum_k \omega_k b_k^\dagger b_k$$

$$\rho_0(t) = e^{-iHt} \rho_0^0 \quad f(t) = \frac{t}{\pi} \int_0^\infty \Gamma(\omega) \frac{\sin^2(\omega t/2)}{\omega^2} \text{odd} \left(\frac{\beta \omega}{2} \right) d\omega$$

$$\Delta E(t) = \frac{2}{\pi} \int_0^\infty \frac{\Gamma(\omega)}{\omega} \sin^2(\frac{\omega t}{2}) d\omega$$

$$\Delta H(t) = \frac{2}{\pi} \int_0^\infty \frac{\Gamma(\omega)}{\omega^2} \sin^2(\frac{\omega t}{2}) d\omega$$

} fällt positiv wenn in des Res. abgelesen

$$\Delta S(t) = S(t) - S(0) + \beta [\Delta E(t) - \lambda \Delta H(t)] \geq 0$$

7.4.3. Period. Treiben

$$\Delta S(t) = \Delta_i S(t) + \sum_{\nu} \beta_{\nu} \Delta Q_{\nu}(t)$$

Langzeitwert sei physikalisch periodisch

$$\text{für große } t^* \quad S(t^* + T) = S(t^*)$$

$$\frac{1}{T} \int_{t^*}^{t^*+T} dt [\dots]$$

Langzeitstrom seien physikalisch periodisch

$$\text{Periodenmittel} \hat{=} \frac{1}{T} \int_{t^*}^{t^*+T} dt [\dots]$$

$$0 = \lim_{t^* \rightarrow \infty} \frac{\Delta_i S(t^*+T) - \Delta_i S(t^*)}{T} + \sum_{\nu} \beta_{\nu} \left[\langle I_{\nu}^{\nu} \rangle - \mu_{\nu} \langle I_{\nu}^{\mu} \rangle \right]$$

$$\underbrace{\Delta_i S(t^*+T)}_{\geq 0} = \underbrace{\Delta_i S(t^*)}_{\geq 0} + \langle G \rangle \cdot T$$

$$\Delta_i S(t^* + n \cdot T) = \Delta_i S(t^*) + \langle G \rangle \cdot n \cdot T \quad n \in \mathbb{Z}^+$$

$$\Rightarrow \langle G \rangle = - \sum_{\nu} \beta_{\nu} \left[\langle I_{\nu}^{\nu} \rangle - \mu_{\nu} \langle I_{\nu}^{\mu} \rangle \right] \geq 0$$

periodengemittelte Ströme (wenn existent) erfüllen 2. HS