

WdH

$$H_I = \sum_x A_x \otimes B_x$$

System
Bad
↑
↑
 A_x^\dagger
 $-B_x^\dagger$

• WdH-Bild bzgl. H_S & H_B

$$\tilde{H}_I = -i \left[\underbrace{e^{+i H_S t} H_I e^{-i H_S t}}_{H_I}, \tilde{P} \right]$$

• Bad-Korrelations-Fkt:

$$C_{\text{op}}(t) = \text{Tr} \left\{ e^{+i H_S t} B_S e^{-i H_S t} B_B \cdot \tilde{\rho}_S \right\} \quad \tilde{\rho}_{\text{red}}(t) = \tilde{\rho}_S(t) \otimes \tilde{\rho}_B$$

• Nicht-Markovsche KGL

$$\dot{\tilde{\rho}}_S = - \sum_{\alpha, \beta} \int_0^t dt' \left\{ C_{\text{op}}(t-t') \left[\tilde{A}_\alpha(t), \tilde{A}_\beta(t') \tilde{\rho}_S(t') \right] + \text{h.c.} \right\} \quad \text{Tr} \tilde{\rho}_S(t) = 1$$

$$= \int_0^t W(t, t') \tilde{\rho}_S(t') dt' \quad \tilde{\rho}_S = \tilde{\rho}_S^\dagger$$

↑
Superoperator

• Markov-Näherung $C_{\text{op}}(t)$ fallen schnell ab

$$\rightarrow \tilde{\rho}_S(t) \rightarrow \tilde{\rho}_S(t)$$

$$\int_0^t \dots dt' \rightarrow \int_0^\infty \dots dt'$$

$$\Rightarrow \dot{\tilde{\rho}}_S = -i [H_S, \tilde{\rho}_S] - \left\{ \sum_{\alpha, \beta} \int_0^\infty C_{\text{op}}(\tau) \left[A_\alpha, e^{-i H_S \tau} A_\beta e^{+i H_S \tau} \tilde{\rho}_S(t) \right] + \text{h.c.} \right\}$$

↑
Korrel. ins Lindblad-Bild

$$= W \tilde{\rho}_S(t) \quad \rightarrow \dot{\tilde{\rho}}_S(t) = e^{W t} \tilde{\rho}_S^0$$

$$\neq \text{Tr} \{ \tilde{\rho}_S \} = 1 \quad \text{h.c.} \quad \tilde{\rho}_S = \tilde{\rho}_S^\dagger \quad \neq \mathbb{1}$$

aber keine Lindblad-Form (in. Allg.)

$$B_S \rho_S B_S^\dagger = \frac{c_1}{2} b_1^\dagger + \frac{c_2}{2} b_2^\dagger + 2 \frac{c_3}{2} b_1^\dagger b_2^\dagger$$

a.) Säkular-Näherung

$$\dot{\tilde{\rho}}_S = \sum_{\omega} \sum_{\omega'} Z_{\omega \omega'} e^{i(\omega-\omega')t} \tilde{\rho}_S(t) = \sum_{\omega} Z_{\omega \omega} \tilde{\rho}_S(t) + \sum_{\omega \neq \omega'} Z_{\omega \omega'} e^{+i(\omega-\omega')t} \tilde{\rho}_S(t)$$

↑
↑
Säkul. Operatoren
Eigenfrequenzen des Systems (Kügelige)
≈ 0
("Säkularnäherung")

$$\dot{\tilde{\rho}}_S = - \sum_{\alpha, \beta} \int_0^\infty C_{\text{op}}(\tau) \left[\tilde{A}_\alpha(t), \tilde{A}_\beta(t-\tau) \tilde{\rho}_S(t) \right] + \text{h.c.} \left. \right\}$$

$$e^{+i k_1 t} A_{\alpha} e^{-i k_2 t} = \sum_{\omega, \omega'} e^{+i E_{\omega} t} |a\rangle \langle a| A_{\alpha} |b\rangle \langle b| e^{-i E_{\omega'} t}$$

$$H_S = \sum_n E_n |a\rangle \langle a| = \sum_{\omega, \omega'} e^{+i(E_{\omega} - E_{\omega'})t} \langle a | A_{\alpha} | b \rangle |a\rangle \langle b|$$

→ Gleichung mit $\delta_{E_{\omega} - E_{\alpha}, E_{\alpha} - E_{\omega}}$ führt zwar auf Lindblad Form

$$\text{Kurzzeitige FT: } \Gamma_{\text{op}}(\omega) = \int_0^\infty C_{\text{op}}(\tau) e^{+i \omega \tau} d\tau = \frac{1}{2} \mathcal{F}_{\text{op}}(\omega) + \frac{1}{2} \mathcal{F}_{\text{op}}^*(\omega)$$

$$\mathcal{F}_{\text{op}}(\omega) = \Gamma_{\text{op}}(\omega) + \Gamma_{\text{op}}^*(\omega) = \int_0^\infty C_{\text{op}}(\tau) e^{+i \omega \tau} d\tau \quad \text{volle FT (gerade)}$$

$$\mathcal{F}_{\text{op}}(\omega) = \Gamma_{\text{op}}(\omega) - \Gamma_{\text{op}}^*(\omega) = \int_0^\infty C_{\text{op}}(\tau) \cdot \text{sgn}(\tau) e^{+i \omega \tau} d\tau \quad \text{volle ungerade FT}$$

$$= \frac{i}{\pi} \oint \frac{\gamma_{\text{res}}(\omega')}{\omega - \omega'} d\omega'$$

Cauchy Hauptwert

Vorr: $H_S = \sum A_n \otimes B_n$ $A_n = A_n^\dagger$ $B_n = B_n^\dagger$

$[\hat{H}_S, \hat{\rho}_S] = 0$ $\text{Tr}\{B_n \hat{\rho}_S\} = 0$ $C_{\text{res}}(\omega) = \text{Tr}\{e^{+i\frac{\omega}{2}\hat{H}_S} B_n e^{-i\frac{\omega}{2}\hat{H}_S} \hat{\rho}_S\}$

beweise $\gamma_{\text{res}}(\omega) = \int C_{\text{res}}(\omega') e^{+i\omega\omega'} d\omega'$ $\hat{G}_{\text{res}}(\omega)$ analog

$\dot{\hat{\rho}}_S = -i[\hat{H}_S + \sum_{ab} \hat{G}_{\text{res}} \hat{L}_{ab}, \hat{\rho}_S] + \sum_{ab, cd} \gamma_{ab, cd} [\hat{L}_{ab} \hat{\rho}_S \hat{L}_{cd}^\dagger - \frac{1}{2}\{\hat{L}_{ab}^\dagger \hat{L}_{ab}, \hat{\rho}_S\}]$

$\hat{L}_{ab} = |a\rangle\langle b|$ $H_S |a\rangle = E_a |a\rangle$

$\hat{G}_{ab} = \sum_{rs} \sum_0 \frac{1}{2i} \hat{G}_{\text{res}}(E_b - E_a) \delta_{E_a, E_b} \langle c|A_b|b\rangle \langle c|A_a|a\rangle^*$

$\gamma_{ab, cd} = \sum_{rs} \gamma_{\text{res}}(E_b - E_a) \delta_{E_b - E_a, E_c - E_a} \langle a|A_b|b\rangle \langle c|A_a|a\rangle^*$

BHS Hermitische Gleichung

+ zeige: $H_S = \sum_{ab} \hat{G}_{ab} \hat{L}_{ab} = H_S^\dagger$

$-\sum_{ab, cd} x_{ab}^* \gamma_{ab, cd} x_{cd} \geq 0$

folgt aus $\gamma_{\text{res}}(\omega)$ ist pos definit

- + L_{ab} sind nicht lokal
- + oft wird H_S vernachlässigt
- + falls $\hat{\rho}_S$ $N \times N \rightarrow \mathcal{L} \hat{\rho}_S = \dot{\hat{\rho}}_S$ $\mathcal{L}: N^2 \times N^2$
- + falls H_S keine Entartungen hat

$\rightarrow \delta_{E_a, E_b} = \delta_{ab}$

$\rightarrow \dot{\hat{\rho}}_{aa} = + \sum_b \gamma_{ab, ab} \hat{\rho}_{bb} - \left(\sum_b \gamma_{ab, ba}\right) \hat{\rho}_{aa}$ Rotationsgleichung

$\gamma_{ab, ab} = \sum_{rs} \gamma_{\text{res}}(E_b - E_a) \langle a|A_b|b\rangle \langle a|A_a|a\rangle^* \geq 0$

$\mathcal{L}_{\text{BAS}} = \begin{pmatrix} \mathcal{L}_{\text{res}} & | & \textcircled{0} \\ \textcircled{0} & | & \mathcal{L}_{\text{coh}} \end{pmatrix}$ $\hat{\rho} = \begin{pmatrix} \hat{\rho}_{\text{res}} \\ \vdots \\ \hat{\rho}_{\text{coh}} \end{pmatrix}$

+ Lindblad Davis - Abb: $\mathcal{L}_{\text{BAS}} e^{-B H_S} = 0$

für $\hat{\rho}_B = \frac{e^{-\beta H_S}}{Z_B}$

1.3.3. Harmon. Oszillator in Stam. Bad

$H_{\text{tot}} = \underbrace{\alpha a^\dagger a}_{H_S} + \underbrace{\sum_k \omega_k b_k^\dagger b_k}_{H_B} + \underbrace{(a + a^\dagger) \sum_k \omega_k (h_k b_k + h_k^\dagger b_k^\dagger)}_{H_I}$ $A = A^\dagger$
 $B = B^\dagger$

$\dot{\hat{\rho}} = -i[\tilde{H}_T, \hat{\rho}]$ $\tilde{H}_T = (a \cdot e^{-i\omega t} + a^\dagger \cdot e^{+i\omega t}) \otimes \sum_k (h_k b_k e^{-i\omega_k t} + h_k^\dagger b_k^\dagger e^{+i\omega_k t})$

$\tilde{a}(t) = e^{+i\omega t a^\dagger} a e^{-i\omega t a}$

$\rightarrow \frac{d}{dt} \tilde{a}(t) = -i\omega \tilde{a}(t)$

$\alpha = \beta = 1$

$$\begin{aligned}
 \text{Tr}\{B \bar{\rho}_0\} &= 0 \\
 C_{11}(\tau) = C(\tau) &= \text{Tr} \left\{ \underbrace{\left(\sum_k b_k b_k^\dagger e^{-i\omega_k \tau} + b_k^\dagger b_k e^{+i\omega_k \tau} \right)}_{\tilde{\gamma}} \left(\sum_k b_k b_k^\dagger + b_k^\dagger b_k \right) \frac{e^{-\beta \sum_k \omega_k b_k^\dagger b_k}}{Z_B} \right\} \\
 &= \frac{1}{2\pi} \int_0^\infty \tilde{\gamma}(\omega) \left[e^{-i\omega \tau} \tilde{\gamma}(1+\omega) + e^{+i\omega \tau} \tilde{\gamma}(\omega) \right] d\omega \\
 \tilde{\gamma}(\omega) &= -\tilde{\gamma}(-\omega) & \tilde{\gamma}(\omega) &= \frac{1}{2\pi} \sum_k \omega_k^2 \delta(\omega - \omega_k) \\
 \tilde{\gamma}(\omega) &= \tilde{\gamma}(\omega) & & \text{Spektrale Dichte} \\
 &= \frac{1}{2\pi} \int \underbrace{\tilde{\gamma}(\omega) \tilde{\gamma}(1+\omega)}_{\gamma_1(\omega)} e^{-i\omega \tau} d\omega
 \end{aligned}$$

$$\dot{\rho} = -i[\rho, a^\dagger a, \rho] - \int_0^\infty C(\tau) \left[(a + a^\dagger), e^{-i\omega a^\dagger a \tau} (a + a^\dagger) e^{+i\omega a^\dagger a \tau} \rho(t) \right] d\tau + \text{h.c.}$$

Karboon-ME (keine Lindblad-Form)

verwandlungsoperator $e^{\pm 2i\omega t}$ im W -Bild

$$\begin{aligned}
 \Gamma(\omega) &= \int_0^\infty C(\tau) e^{+i\omega \tau} d\tau \Rightarrow \frac{1}{2} \gamma + \frac{i}{2} \beta = \Gamma(+\omega) \\
 \frac{1}{2} \tilde{\gamma} + \frac{i}{2} \beta &= \Gamma(-\omega)
 \end{aligned}$$

$$\Rightarrow \dot{\rho} = -i \left[\frac{\beta}{2} a^\dagger a + \frac{\beta}{2} a a^\dagger, \rho \right] + \gamma \left[a^\dagger a \rho a^\dagger - \frac{1}{2} \{a^\dagger a, \rho\} \right] + \tilde{\gamma} \left[a^\dagger \rho a - \frac{1}{2} \{a a^\dagger, \rho\} \right]$$

$$\Rightarrow \dot{\rho} = -i \left[\omega a^\dagger a + \left(\frac{\beta}{2} + \frac{\beta}{2} \right) a^\dagger a, \rho \right] + \tilde{\gamma}(\omega) \tilde{\gamma}(1+\omega) \left[a \rho a^\dagger - \frac{1}{2} \{a^\dagger a, \rho\} \right] + \tilde{\gamma}(\omega) \tilde{\gamma}(\omega) \left[a^\dagger \rho a - \frac{1}{2} \{a a^\dagger, \rho\} \right]$$

$$H_0(a) = \omega a^\dagger a \Rightarrow E_n(a)$$

$$a^\dagger a = \sum_{k=1}^{\infty} k \cdot |k\rangle \langle k|$$

$$a = \sum_{k=1}^{\infty} \sqrt{k} |k-1\rangle \langle k|$$

1.3.4. Gleichgewichts-TD

Wann ist $\sum e^{-\beta \omega_k} = 0$

falls $\rho_0 = \frac{e^{-\beta H_0}}{Z_B}$

$C_{\beta, \tau}(\bar{v}) = C_{\beta, \tau}(-\bar{v} - i\beta)$ Kubo-Martin-Schwinger Relation

$$\begin{aligned}
 C_{\beta, \tau}(-\bar{v} - i\beta) &= \text{Tr} \left\{ \underbrace{e^{+i\omega_0(-\bar{v} - i\beta)} B_p e^{-i\omega_0(-\bar{v} - i\beta)} B_r \frac{e^{-\beta H_0}}{Z_B}}_{\text{KMS}} \right\} \\
 &= \text{Tr} \left\{ \underbrace{e^{-\beta H_0} e^{+i\omega_0 \bar{v}} B_r \frac{e^{-\beta H_0}}{Z_B} e^{+i\omega_0 \bar{v}} B_p}_{\text{KMS}} \right\} = \text{Tr} \left\{ B_r(\tau) B_p \frac{e^{-\beta H_0}}{Z_B} \right\} \\
 &= C_{\beta, \tau}(\bar{v})
 \end{aligned}$$

$$\frac{\omega_0(\omega)}{1 + \omega_0(\omega)} = e^{-\beta \omega}$$

Korrig.: 2. Kompakt