

WdK
WZ-Verteilung

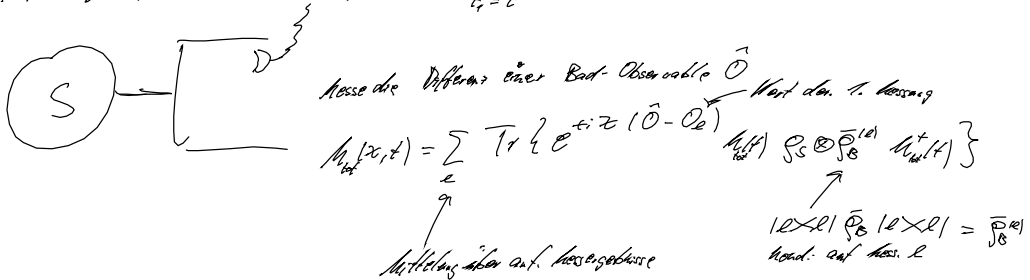
$$W^i(t) = \frac{\text{Tr} \{ Z_i \cdot e^{z_i \cdot \hat{O}} \cdot Z_i^\dagger \}}{\text{Tr} \{ Z_i \cdot \hat{\rho} \}}$$

\uparrow 2. Sprung \uparrow 1. Sprung
 \uparrow kond. auf 1. Sprung \uparrow langzeit-WZ-Verteilung

mittlerer WZ zu Sprüngen (j) und (\bar{j})

$$k\text{-tes Moment } \langle \hat{O}^k \rangle = \int_0^\infty W^i(t) \cdot \hat{O}^k dt$$

• Mikroskop Ableitung von ZF 2-Messungen Bed: $t_0=0$
 $t_1=t$



z.B. $(-i \hat{O}_z)^k M_{\text{det}}(z,t) |_{z=0} = \sum_P \text{Tr} \{ (\hat{O} - O_0)^k \rho_{\text{stat}}(t) \}$

→ verallg. Zeit-Evol.-Op. $M_{z_1, z_2}(t) = e^{-i \hat{O} z_1 t} U(t) e^{-i \hat{O} z_2 t} \sum_P \bar{\rho}_B^{(1)}$

→ Abl. der Mastergl. z.B.: $[M_z + t \cdot \lambda(z)] \dots \rho_S^0 \stackrel{!}{=} \text{Tr}_B \{ M_{z_1, z_2}(t) \rho_S^0 \otimes \bar{\rho}_B M_{z_1, z_2}^+(t) \}$

→ in der Mastergleichung bekommt die Streifen ein Zeitfeld "L_z ∈ L_B⁺ → L_z ∈ L_B⁺ · e^{t·z}"

→ verallg. Korrelationsfunktion

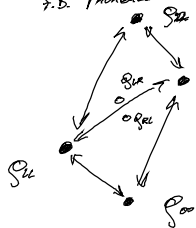
$$C_{\text{op}}(t) \rightarrow C_{\text{op}}^z(t) = \text{Tr}_B \{ e^{-i \hat{O} x} e^{+i t_0 \hat{P}} \rho_B e^{-i t_0 \hat{O}} e^{+i \hat{O} x} \rho_B \cdot \bar{\rho}_B \}$$

falls $\bar{\rho}_B = \sum_P P_e 12 \times 12 \Leftrightarrow [\bar{\rho}_B, \hat{O}] = 0$

→ $\bar{\rho}_B = \rho_B$

Wann sinnvoll

+ Kohärenzen
z.B. Phänomene-assul. Tunnel



Treatment mit mikroskop. ZF möglich

3.4.1. Beispiel: SRL $\hat{O} = H_B$

$$H = \underbrace{\varepsilon d^\dagger d}_{H_S} + d \otimes \sum_k \underbrace{t_k}_{A_1} c_k^\dagger + d \otimes \sum_k \underbrace{t_k^*}_{A_2} c_k + \sum_k \underbrace{\varepsilon_k}_{H_B} c_k^\dagger c_k$$

$$\begin{aligned} e^{+i\alpha c_k^\dagger c_k} c_k^\dagger e^{-i\alpha c_k^\dagger c_k} \\ = c_k^\dagger \cdot e^{+i\alpha} \end{aligned}$$

$$\begin{aligned} C_{12}^z(\tau) &= \sum_{k_1, k_2} \text{Tr}_B \left\{ t_{k_1} c_{k_1}^\dagger e^{+i\varepsilon_{k_1}\tau} e^{-i\varepsilon_{k_2}\tau} t_{k_2}^* c_{k_2} \bar{\rho}_B \right\} \\ &= \sum_k (t_k)^2 \cdot f(\varepsilon_k) \cdot e^{+i\varepsilon_k(\tau-\tau)} \\ &= \frac{1}{2\pi} \int \Gamma(\omega) f(\omega) e^{+i\omega\tau} e^{-i\omega\tau} d\omega \\ &= \frac{1}{2\pi} \int \underbrace{\Gamma(\omega) \cdot f(\omega)}_{\gamma_{12}^z(\omega)} e^{+i\omega\tau} e^{-i\omega\tau} d\omega \end{aligned}$$

analog: $C_{21}^z(\tau) = \frac{1}{2\pi} \int \underbrace{\Gamma(\omega) [1-f(\omega)]}_{\gamma_{21}^z(\omega)} e^{+i\omega\tau} e^{-i\omega\tau} d\omega$

→ BHS Perturbations (oder GG mit $\tau \rightarrow \infty$)

$$\gamma_{ab,ab}^z = \sum_{\gamma_B} \gamma_{\gamma_B}^z (E_B - E_a) \langle a | A_{\gamma_B} | b \rangle \langle a | A_{\gamma_B}^\dagger | b \rangle$$

$$E_0 = 0 \quad E_1 = \varepsilon$$

$$\begin{aligned} \gamma_{01,01}^z &= \gamma_{11}^z(+\varepsilon) = \Gamma(\varepsilon) [1-f(\varepsilon)] e^{+i\varepsilon\tau} \\ \gamma_{10,10}^z &= \gamma_{11}^z(-\varepsilon) = \Gamma(\varepsilon) f(\varepsilon) \cdot e^{-i\varepsilon\tau} \end{aligned}$$

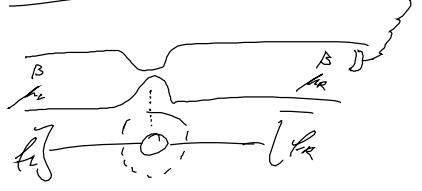
Konvention: positiv, wenn in der Rec. gerichtet

$$\hookrightarrow Z(z) = \Gamma(\varepsilon) \begin{pmatrix} -f(\varepsilon) & + [1-f(\varepsilon)] e^{+i\varepsilon\tau} \\ f(\varepsilon) \cdot e^{-i\varepsilon\tau} & - [1-f(\varepsilon)] \end{pmatrix}$$

→ konsistent

analog $\hat{O} = H_B$

3.4.2. Beispiel QPC $\hat{O} = H_{B,QPC}$



$$H_{SET} = \varepsilon d^\dagger d + \sum_{\substack{a,v \\ \text{verteilt}}} (t_{av} d c_{av}^\dagger + t_{av}^* c_{av}^\dagger d) + \sum_{k,v} \varepsilon_{kv} c_{kv}^\dagger c_{kv}$$

$$H_{QPC} = \sum_k \sum_{\text{verteilt}} \varepsilon_{kv} \gamma_{kv}^+ \gamma_{kv}$$

$$H_{\text{SET-QPC}} = \underbrace{\left(\mathbb{1} - \tilde{v} \cdot \text{d.t.d.} \right)}_A \sum_{kk'} \underbrace{\left(t_{kk'} \gamma_{kL} \gamma_{kR}^+ + t_{kk'}^* \gamma_{kR} \gamma_{kL}^+ \right)}_B$$

Reduktion des Tunnels bei besetzten SET-QD
 $0 < \tilde{v} < 2$

betrachte nur die Kopplung an das QPC

$$\hat{D} = \sum_n \gamma_{nL}^+ \gamma_{nR}$$

$$\mathcal{P}_{B, \text{QPC}} = \frac{e^{-\beta \sum_k (E_{kL} - \mu_L) \gamma_{kL}^+ \gamma_{kL}}}{Z_L} \frac{e^{-\beta \sum_k (E_{kR} - \mu_R) \gamma_{kR}^+ \gamma_{kR}}}{Z_R}$$

$$C^{\tilde{v}}(\tilde{\omega}) = \text{Tr} \left\{ e^{-i\tilde{v}\tilde{\omega}\hat{D}} e^{+i\tilde{v}\text{d.t.d.}\tilde{v}} B e^{-i\tilde{v}\text{d.t.d.}\tilde{v}} e^{+i\tilde{v}\tilde{\omega}\hat{D}} B \bar{\mathcal{P}}_B \right\} \quad (\bar{\mathcal{P}}_B = \bar{\mathcal{P}}_B)$$

$$= \sum_{kk'} \text{Tr} \left\{ \left[t_{kk'} \gamma_{kL} \gamma_{kR}^+ e^{+i(E_{kR} - E_{kL})\tilde{v}} e^{-i\tilde{v}\tilde{\omega}} + t_{kk'}^* \gamma_{kR} \gamma_{kL}^+ e^{-i(E_{kR} - E_{kL})\tilde{v}} e^{+i\tilde{v}\tilde{\omega}} \right] \right.$$

$$\left. \times \left[t_{kk'} \gamma_{kL} \gamma_{kR}^+ + t_{kk'}^* \gamma_{kR} \gamma_{kL}^+ \right] \mathcal{P}_{B, \text{QPC}} \right\}$$

$$= \sum_{kk'} \left\{ e^{-i\tilde{v}\tilde{\omega}} |t_{kk'}|^2 e^{+i(E_{kR} - E_{kL})\tilde{v}} [1 - f_L(E_{kL})] \cdot f_R(E_{kR}) \right.$$

$$\left. + e^{+i\tilde{v}\tilde{\omega}} |t_{kk'}|^2 e^{-i(E_{kR} - E_{kL})\tilde{v}} f_L(E_{kL}) [1 - f_R(E_{kR})] \right\}$$

$$T(\omega, \omega') = 2\pi \sum_{kk'} |t_{kk'}|^2 \delta(\omega - E_{kL}) \delta(\omega' - E_{kR}) \quad \text{Transmission}$$

$$\gamma(\omega) = \int C(\tilde{\omega}) e^{+i\tilde{\omega}\omega} d\tilde{\omega}$$

$$\rightarrow = e^{-i\tilde{v}\tilde{\omega}} \frac{1}{2\pi} \int d\omega \int d\omega' T(\omega, \omega') e^{+i(\omega - \omega')\tilde{\omega}} [1 - f_L(\omega)] f_R(\omega')$$

$$+ e^{+i\tilde{v}\tilde{\omega}} \frac{1}{2\pi} \int d\omega \int d\omega' T(\omega, \omega') e^{-i(\omega - \omega')\tilde{\omega}} f_L(\omega) [1 - f_R(\omega')] \quad \frac{1}{2\pi} \int e^{i\tilde{\omega}\omega} d\tilde{\omega} \delta(\tilde{\omega})$$

$$\gamma^{\tilde{v}}(\omega) = e^{-i\tilde{v}\tilde{\omega}} \int d\omega T(\omega, \omega - \omega) [1 - f_L(\omega)] f_R(\omega - \omega) \quad \text{Transfer von } R \rightarrow L \text{ unter Absorption von Energie } \omega \text{ aus der Spalte}$$

$$+ e^{+i\tilde{v}\tilde{\omega}} \int d\omega T(\omega, \omega + \omega) f_L(\omega) [1 - f_R(\omega + \omega)] \quad \text{Transfer von } L \rightarrow R$$

Integrale des Typs $I = \int d\omega f_L(\omega) [1 - f_R(\omega)]$

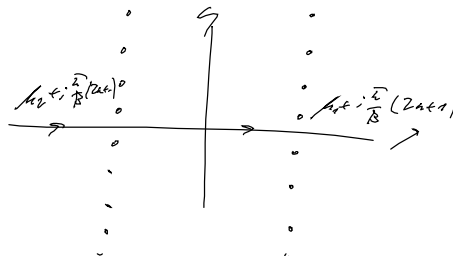
Vereinfachung $T(\omega, \omega') \approx T$

a) für $\beta \rightarrow \infty$ (Temp. $\rightarrow 0$)

$$I \rightarrow (f_L - f_R) \cdot \Theta(\mu_L - \mu_R)$$

b) für endl. Temp.

Funktionentheorie



$$\dots T = \frac{\Gamma_L \Gamma_R}{1 - e^{-\beta(\mu_L - \mu_R)}}$$

$$\Gamma_{a,b}^{\pm} = \Gamma^{\pm} (E_a - E_b) |\langle a | [\hat{H} - \tilde{V} d^{\dagger} d] | b \rangle|^2$$

$$\Sigma_{QPC}(\xi) = \begin{pmatrix} 1 & 0 \\ 0 & (1-\alpha)^2 \end{pmatrix}^T \left[(e^{-\beta\xi} - 1) \frac{V}{e^{\beta V} - 1} + (e^{\beta\xi} - 1) \frac{V}{1 - e^{-\beta V}} \right]$$

$$V = \mu_L^{QPC} - \mu_R^{QPC}$$

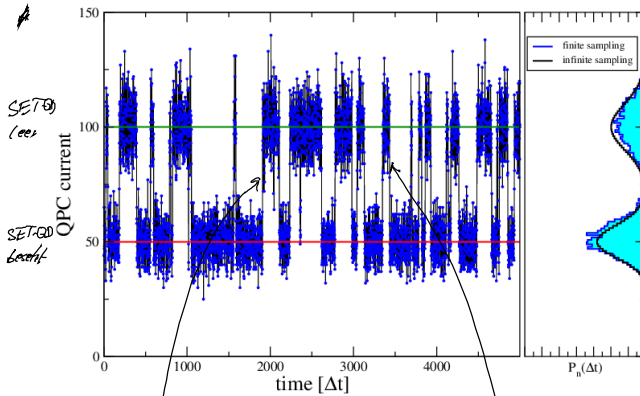
$$\Sigma_{ges}(\xi) = \Sigma_{SET} + \Sigma_{QPC}(\xi)$$

a.) z.B. $\Gamma_L = \Gamma_R \rightarrow 0$

$$\begin{matrix} \uparrow & & \uparrow \\ I_0 = T \cdot V & & I_1 = (1-\alpha)^2 \cdot I_0 \end{matrix}$$

SET-QD leer falls: $\{\Gamma_L, \Gamma_R\} \ll \{T \cdot V, (1-\alpha)^2 \cdot T \cdot V\}$

b) $P_{kur}(dt) = \frac{1}{2\pi} \int Tr \{ e^{i \int_{t_0}^{t_0+dt} \xi(t) dt - i \lambda \xi} \} d\xi$



$$\bar{I}_{QPC} = \frac{\Sigma R}{\Delta \phi}$$

SET entleert sich (ins rechte Bad)

SET füllt sich (aus dem linken Bad)